

Quasi-ergodic theorems for Feynman-Kac semigroups and large deviation for additive functionals

Takara Tagawa*

(joint work with D. Kim and A. Velleret)

1 Introduction

This talk is based on [2, 3]. Let E be a locally compact separable metric space and $\mathbf{m}$ be a positive Radon measure on E with full topological support. Let $X = (\Omega, \mathcal{M}, \mathcal{M}_t, X_t, \mathbb{P}_x, \zeta)$ be a strong Markov process on E with the lifetime ζ . We use $\mathbb{E}_x$ to denote the expectation with respect to $\mathbb{P}_x$ for any $x \in E$. Denote by $\{p_t\}_{t \geq 0}$ and $\{R_\alpha\}_{\alpha \geq 0}$ the transition semigroup and resolvent of X , respectively.

Insightful properties of an almost surely killed Markov process X (i.e., $\mathbb{P}_x(\zeta < \infty) = 1$ for any $x \in E$) when conditioned to survive have been derived by the study of the long-time behavior of the mean-ratio of a continuous additive functional. In [1, 4], the authors established the following quasi-ergodic limit theorem: there exists a (probability) measure η on E such that for any $f \in L^1(E; \eta)$ and $x \in E$,

$$(1.1) \quad \lim_{t \rightarrow \infty} \mathbb{E}_x \left[\frac{1}{t} \int_0^t f(X_s) ds \mid t < \zeta \right] = \int_E f(x) \eta(dx).$$

The measure η is often called a quasi-ergodic distribution of X . Let G be a measurable function on $E \times E$ vanishing on the diagonal. Define a purely discontinuous additive functional of X by $A_t^G := \sum_{0 < s \leq t} G(X_{s-}, X_s)$. This additive functional often appears when considering the pure jump effects in Markov processes and, in a particular case, is thought to express the number of jumps in pure jump processes. In [3], we propose to replace the integration in (1.1) over the full path of the processes by a summation over all jumps of the process up to time t (by taking into account both the position before and the position after the jump) and established a quasi-ergodic theorem for moments of the mean-ratio of A_t^G caused by the pure jump effects of symmetric Markov processes.

Let μ be a positive smooth measure on E whose associated positive continuous additive functional of X is denoted by A_t^μ . Let F be a positive symmetric bounded measurable function on $E \times E$ vanishing on the diagonal, and put $A_t^{\mu, F} := A_t^\mu + A_t^F$. In combination with the extinction that plays a similar role, we are led to consider the following family of non-local Feynman-Kac transforms, indexed by a time t which can be thought of as a final time for the biased action: for $x \in E$ and $\Lambda \in \mathcal{M}$,

$$(1.2) \quad \mathbb{P}_{x:t}^{\mu, F}(\Lambda) := \mathbb{P}_x \left(\Lambda \cdot W_t^{\mu, F} \right) \quad \text{where} \quad W_t^{\mu, F} := \exp \left(A_t^{\mu, F} \right) \mathbf{1}_{\{t < \zeta\}}.$$

With (1.2), we define a probability measure by $\mathbb{P}_{x|t}^{\mu, F}(\Lambda) := \mathbb{P}_{x:t}^{\mu, F}(\Lambda) / \mathbb{P}_{x:t}^{\mu, F}(\Omega)$ for $x \in E$ and $\Lambda \in \mathcal{M}$. We use $\mathbb{E}_{x|t}^{\mu, F}$ to denote the expectation with respect to $\mathbb{P}_{x|t}^{\mu, F}$.

The main subject of this talk is to discuss a quasi-ergodic limit theorem for the additive functional

$$A_t^{f, G} := \int_0^t f(X_s) ds + \sum_{0 < s \leq t} G(X_{s-}, X_s)$$

under the symmetric Markov process X driven by the *non-local Feynman-Kac transform* (1.2).

*Department of Mathematics and Engineering, Graduate School of Science and Technology, Kumamoto University

2 Main results

Let X is an $\mathbf{m}$ -symmetric irreducible explosive Markov process on E satisfying the strong Feller property (that is, $p_t(\mathfrak{B}_b(E)) \subset C_b(E)$ for any $t > 0$) and the tightness property (that is, for any $\varepsilon > 0$ there exists a compact $K \subset E$ such that $\sup_{x \in E} R_1 1_{K^c}(x) \leq \varepsilon$). Let $(N(x, dy), H_t)$ be a Lévy system for X . We assume that $\mu_H = \mathbf{m}$ (so that $H_t = t$ for $t < \zeta$), where μ_H is the Revuz measure of the positive continuous additive functional H_t .

A positive smooth measure μ on E (resp. a positive measurable function F on $E \times E$) is said to be in the Kato class associated to X , if it satisfies

$$\limsup_{t \rightarrow 0} \sup_{x \in E} \mathbb{E}_x [A_t^\mu] = 0 \quad \left(\text{resp.} \quad \limsup_{t \rightarrow 0} \sup_{x \in E} \mathbb{E}_x \left[\int_0^t \int_E F(X_s, y) N(X_s, dy) ds \right] = 0 \right).$$

Under the Kato class conditions on μ and F , there is the ground state $\phi_0^{(\theta)} := \phi_0^{\theta\mu, \theta F}$ of the Feynman-Kac semigroup $p_t^{\theta\mu, \theta F} f(x) := \mathbb{E}_x[W_t^{\theta\mu, \theta F} f(X_t)]$, $\theta \in \mathbb{R}$. We say that $p_t^{\theta\mu, \theta F}$ is θ -intrinsically ultracontractive (**(IUC)** $_\theta$ in short), if for any $\theta \in \mathbb{R}$ there exists a constant $c_t(\theta) > 0$ such that $p_t^{\theta\mu, \theta F}(x, y) \leq c_t(\theta) \phi_0^{(\theta)}(x) \phi_0^{(\theta)}(y)$ for all $t > 0$ and $x, y \in E$. Many examples are satisfying **(IUC)** $_\theta$ (cf. [2, 3]). Write $\phi_0 := \phi_0^{(1)}$. Define a measure $\mathcal{J}_{\phi_0}$ on $E \times E$ by

$$\mathcal{J}_{\phi_0}(dxdy) := \phi_0(x) \phi_0(y) \exp(F(x, y)) N(x, dy) \mathbf{m}(dx)$$

One of our main results is as follows:

Theorem 2.1. *Let μ and F be of the Kato class one associated to X . Assume **(IUC)** $_1$. Then, for any $f \in L^1(E; \phi_0^2 \mathbf{m})$, $G \in \mathfrak{B}(E \times E \setminus \text{diag})$ satisfying $N^{\phi_0}[G] \in L^1(E; \phi_0^2 \mathbf{m})$ and any $x \in E$,*

$$\lim_{t \rightarrow \infty} \mathbb{E}_{x|t}^{\mu, F} \left[\frac{1}{t} A_t^{f, G} \right] = \int_E f \phi_0^2 d\mathbf{m} + \iint_{E \times E} G(x, y) \mathcal{J}_{\phi_0}(dxdy).$$

For a positive $V \in L^1(E; \phi_0^2 \mathbf{m})$, let $(\mathcal{E}^{V, F}, \mathcal{D}(\mathcal{E}))$ be the quadratic form given by

$$\mathcal{E}^{V, F}(u, v) := \mathcal{E}(u, v) - \int_E u(x) v(x) V(x) \mathbf{m}(dx) - \iint_{E \times E} u(x) v(y) (\exp(F(x, y)) - 1) N(x, dy) \mathbf{m}(dx),$$

where $(\mathcal{E}, \mathcal{D}(\mathcal{E}))$ is the Dirichlet form associated with X . For $\theta > 0$, define the spectral function

$$C_{V, F}(\theta) := -\inf \left\{ \mathcal{E}^{\theta V, \theta F}(u, u) \mid u \in \mathcal{D}(\mathcal{E}), \int_E u^2 d\mathbf{m} = 1 \right\}.$$

By the analytic perturbation theory, $C_{V, F}(\theta)$ is differentiable in θ and its derivative $C'_{V, F}(\theta)$ is strictly increasing. Put $\Psi_{V, F}(\theta) := C'_{V, F}(\theta)$. As an application of Theorem 2.1, we have the following LDP for $A_t^{V, F}$:

Theorem 2.2. *Let V and F be of the Kato class one associated to X . Assume **(IUC)** $_\theta$ for any $\theta \in \mathbb{R}$. Then, for any $\gamma \in \Psi_{V, F}(\mathbb{R}_+)^o := \{\Psi_{V, F}(\theta) : \theta \in \mathbb{R}_+\}^o$ and any $x \in E$,*

$$\lim_{t \rightarrow \infty} \frac{1}{t} \log \mathbb{P}_x \left(\frac{1}{t} A_t^{V, F} \in [\gamma, \infty), t < \zeta \right) = C_{V, F}(\theta_\gamma) - \theta_\gamma C'_{V, F}(\theta_\gamma), \quad \text{where } \theta_\gamma = \Psi_{V, F}^{-1}(\gamma).$$

References

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