

The spectrum of non-symmetric operators and Markov processes

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1. Normal operators

General framework

- H : a **complex** Hilbert space
- T : a closed operator with domain $\text{Dom}(T)$
- $\Theta(T)$: the numerical range of T defined by

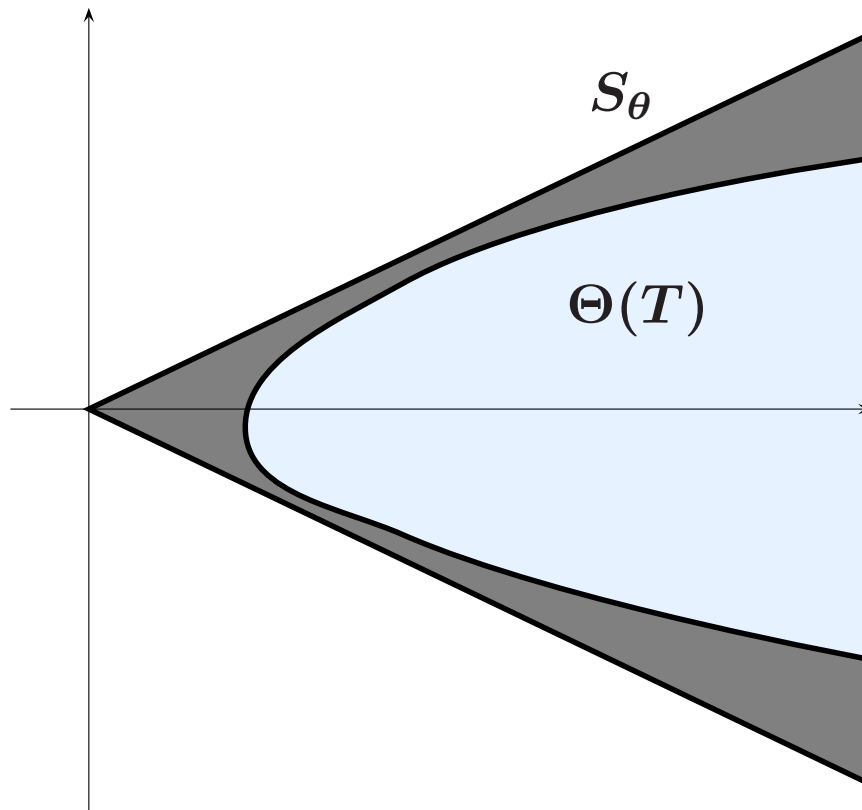
$$\Theta(T) := \{(Tu, u); u \in \text{Dom}(T)\}.$$

- T is called **accretive** if

$$\Re(Tu, u) \geq 0, \quad \forall u \in \text{Dom}(T)$$

- T is called **m -accretive** if $\text{Ran}(T - \zeta) = H$ for some $\zeta \in \mathbb{C}$.

- T is called **sectorial** if $\Theta(T) \subseteq S_\theta$, $\theta \in [0, \frac{\pi}{2})$ where $S_\theta = \{z \in \mathbb{C}; |\arg z| \leq \theta\}$.



- T is called **quasi-sectorial** if $T + \gamma$ is sectorial for some $\gamma > 0$.

Normal operators

- A is called **normal** if

$$A^* A = A A^*$$

- A has an spectral decomposition:

$$A = \int_{\mathbb{C}} z E(dz)$$

- A^* : an adjoint operator of A .

$$A^* = \int_{\mathbb{C}} \bar{z} E(dz)$$

- $\overline{\Theta(A)} = \overline{\text{co}(\sigma(A))}$

From now on, we assume that A is normal and m -accretive.

- \sqrt{A} is defined by

$$\sqrt{A} = \int_{\mathbb{C}} \sqrt{z} E(dz)$$

with

$$\text{Dom}(\sqrt{A}) = \left\{ u \in H; \int_{\mathbb{C}} |z| (u, E(dz)u) < \infty \right\}.$$

- $\text{Dom}(\sqrt{A}) = \text{Dom}(\sqrt{A^*})$
- a : a sesquilinear form associated with A is given by

$$a(u, v) = (Au, v), \quad u, v \in \text{Dom}(A)$$

- A symmetric part of a is defined by

$$b(u, v) = \frac{(Au, v) + (A^*u, v)}{2}, \quad u, v \in \text{Dom}(A).$$

- b can be written

$$b(u, v) = \int_{\mathbb{C}} \Re z(u, E(dz)v).$$

- $(b, \text{Dom}(b))$ is closed where

$$\text{Dom}(b) = \{u \in H; \int_{\mathbb{C}} \Re z(u, E(dz)u) < \infty\}.$$

- $\text{Dom}(\sqrt{A}) \subseteq \text{Dom}(b)$

Theorem 1. $\text{Dom}(\sqrt{A}) = \text{Dom}(b)$ if and only if $1 + \sigma(A) \subseteq S_{\theta}$ for some $\theta \in (0, \pi/2)$.

2. Normal operators and generalized Dirichlet forms

Stannat (1994) introduced the generalized Dirichlet form.

We will show that Markovian semigroup generated by a normal operator can be formulated in the framework of generalized Dirichlet form.

- M : a Hausdorff topological space
- (M, m) : σ -finite measure space
- $H = L^2(m)$
- \mathfrak{A} : a normal operator
- We assume that \mathfrak{A} and \mathfrak{A}^* is m -dissipative (i.e., $-\mathfrak{A}$ and $-\mathfrak{A}^*$ is m -accretive)

By spectral decomposition,

$$(1) \quad -\mathfrak{A} = \int_{\mathbb{C}} z E(dz).$$

We define

$$(2) \quad -L = \int_{\mathbb{C}} \Re z E(dz), \quad -\Lambda = \int_{\mathbb{C}} i \Im z E(dz).$$

L and $i\Lambda$ are self-adjoint with domains

$$\text{Dom}(L) = \left\{ f; \int_{\mathbb{C}} |\Re z|^2 (f, E(dz) f) < \infty \right\},$$

$$\text{Dom}(\Lambda) = \left\{ f; \int_{\mathbb{C}} |\Im z|^2 (f, E(dz) f) < \infty \right\}.$$

L generates a semigroup. Symmetric bilinear form $\tilde{\mathcal{E}}$ is defined by

$$\tilde{\mathcal{E}}(f, g) = \int_{\mathbb{C}} \Re z (f, E(dz) g)$$

with the domain

$$\text{Dom}(\tilde{\mathcal{E}}) = \left\{ f; \int_{\mathbb{C}} |\Re z| (f, E(dz) f) < \infty \right\}.$$

We set $\mathcal{V} = \text{Dom}(\tilde{\mathcal{E}})$.

Similarly, Λ generates a semigroup denoted by $\{U_t\}_{t \geq 0}$.

Proposition 2. $\{U_t\}$ is a C_0 -semigroup in \mathcal{V} .

We regard $\Lambda: \text{Dom}(\Lambda) \cap \mathcal{V} \rightarrow \mathcal{V}'$ as an operator from \mathcal{V} to \mathcal{V}' . Its closure is denoted by (Λ, \mathcal{F}) .

Proposition 3. $f \in \mathcal{F}$ if and only if

$$\int_{\mathbb{C}} \left(\frac{|\Im z|^2}{\Re z + 1} + \Re z \right) (f, E(dz)f) < \infty.$$

Similar argument can be done for the dual semigroup \hat{U}_t of U_t . The generator is

$$\hat{\Lambda} = - \int_{\mathbb{C}} i \Im z E(dz).$$

Now we can apply the theory of generalized Dirichlet form. The Dirichlet form is defined by

$$\mathcal{E}(f, g) = \begin{cases} \tilde{\mathcal{E}}(f, g) - \langle \Lambda f, g \rangle, & \text{if } f \in \mathcal{F}, g \in \mathcal{V}, \\ \tilde{\mathcal{E}}(f, g) - \langle \hat{\Lambda} g, f \rangle, & \text{if } f \in \mathcal{V}, g \in \hat{\mathcal{F}}. \end{cases}$$

Assuming the Markovian property, we can define the capacity.

We assume the **quasi-regularity** of \mathcal{E} . Now applying the following theorem, we can get a Markov process associated with \mathfrak{A} .

Theorem 4. (Stannat 1994) Under the following condition (D3), there exists an **m -thght special standard process**.

(D3) There exists a linear subspace $\mathcal{Y} \subseteq L^2(m) \cap L^\infty(m)$ such that $\mathcal{Y} \cap \mathcal{F}$ is dense in \mathcal{F} , $\lim_{\alpha \rightarrow \infty} e_{\alpha G_\alpha} u - u = 0$ in H for all $u \in \mathcal{Y}$ and for the closure $\overline{\mathcal{Y}}$ of \mathcal{Y} in $L^\infty(m)$ it follows that $u \wedge \alpha \in \overline{\mathcal{Y}}$ for $u \in \mathcal{Y}$ and $\alpha \geq 0$.

3. Criterion for normal operators

- H : a Hilbert space
- A, B : accretive operators on \mathcal{D}
- Assume that $\overline{A}, \overline{B}$ are m -accretive

Theorem 5. Assume that $A\mathcal{D} \subseteq \mathcal{D}, B\mathcal{D} \subseteq \mathcal{D}$ and

$$AB = BA \quad \text{on } \mathcal{D},$$

$$(Au, v) = (u, Bv), \quad u, v \in \mathcal{D}.$$

Then \overline{A} is normal and $\overline{A}^* = \overline{B}$.

Examples on a Riemannian manifold

- M : a complete Riemannian manifold
- m : the Riemannian volume

We take a function $U \in C^\infty(M)$ and define a measure ν by

$$\nu = e^{-U} m$$

Define an operator on $H = L^2(\nu)$ by

$$\mathfrak{A} = \frac{1}{2} \Delta_\nu + b$$

where $\Delta_\nu = -\nabla_\nu^* \nabla$. Then

$$\mathfrak{A}_\nu^* = \frac{1}{2} \Delta_\nu - b - \operatorname{div}_\nu b.$$

Here div_ν denotes the divergence with respect to ν .

We give a criterion for $\mathfrak{A} = \Delta_\nu + b$ being a normal operator.

Theorem 6. Assume that $\operatorname{div}_\nu b$ is bounded from below. Then \mathfrak{A} is normal if and only if b is a **Killing vector field** and the following identities hold:

$$\begin{aligned} \left(\frac{1}{2}\Delta_\nu + b\right) \operatorname{div}_\nu b &= 0, \\ [\nabla U, b] + \nabla \operatorname{div}_\nu b &= 0. \end{aligned}$$

4. One-dimensional Brownian motion with a drift

We consider an operator $\mathfrak{A} = \frac{d^2}{dx^2} - c \frac{d}{dx}$ on $L^2(\mathbb{R}, \nu_1)$. Here ν_1 is a measure defined by

$$(3) \quad \nu_1(dx) = e^{-cx} dx.$$

Then \mathfrak{A} is a self-adjoint operator with

$$(\mathfrak{A}f, g) = - \int_{\mathbb{R}} f'(x)g'(x) \nu_1(dx).$$

To investigate the spectrum of \mathfrak{A} , we use the following isometric map $I: L^2(\nu_1) \longrightarrow L^2(dx)$:

$$If(x) = e^{-cx/2} f(x).$$

We have

$$I \circ \mathfrak{A} \circ I^{-1} = \frac{d^2 f}{dx^2} - \frac{c^2}{4},$$

i.e., the following diagram is commutative:

$$\begin{array}{ccc}
 L^2(\nu_1) & \xrightarrow{\mathfrak{A}} & L^2(\nu_1) \\
 I \downarrow & & \downarrow I \\
 L^2(dx) & \xrightarrow{\frac{d^2}{dx^2} - \frac{c^2}{4}} & L^2(dx)
 \end{array}$$

Hence the spectrum $-\mathfrak{A}$ is

$$(4) \quad \sigma(-\mathfrak{A}) = \left[\frac{c^2}{4}, \infty \right).$$

We now consider an perturbation of \mathfrak{A} . Let b be an vector field defined by

$$b = k \frac{d}{dx}.$$

We consider an operator of the form $\mathfrak{A} + b$. We are interested in how the spectrum changes. b is clearly an Killing vector field. The divergence of b with respect to ν_1

$$\operatorname{div}_{\nu_1} b = -ck$$

and so it satisfies

$$\begin{aligned}(\mathfrak{A} + b) \operatorname{div}_\nu b &= 0, \\ [(\nabla U)^\sharp, b] + \nabla \operatorname{div}_\nu b &= 0.\end{aligned}$$

Here $U(x) = cx$. By Theorem 6, $\mathfrak{A} + b$ is a normal operator. Under the transformation of I , we have

$$I \circ (\mathfrak{A} + b) \circ I^{-1} = \frac{d^2}{dx^2} + k \frac{d}{dx} - \frac{c(c - 2k)}{4}.$$

It is enough to get the spectrum of $\frac{d^2}{dx^2} + k \frac{d}{dx}$. Recall the Fourier transform as

$$\hat{f}(\xi) = \frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} f(x) e^{-i\xi x} dx.$$

This gives an isometry from $L^2(dx)$ onto $L^2(d\xi)$. Note that

$$\int_{\mathbb{R}} \left(\frac{d^2}{dx^2} + k \frac{d}{dx} \right) f(x) \overline{g(x)} dx = \int_{\mathbb{R}} (-\xi^2 + ik\xi) \hat{f}(\xi) \overline{\hat{g}(\xi)} d\xi$$

which means that

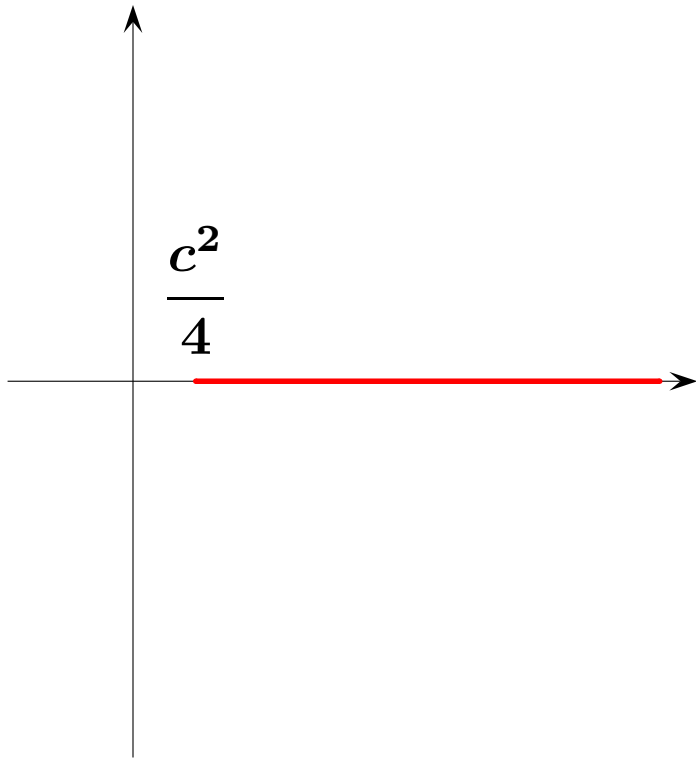
$$\sigma\left(\frac{d^2}{dx^2} + k\frac{d}{dx}\right) = \{-\xi^2 + ik\xi; \xi \in \mathbb{R}\}.$$

Theorem 7. We have

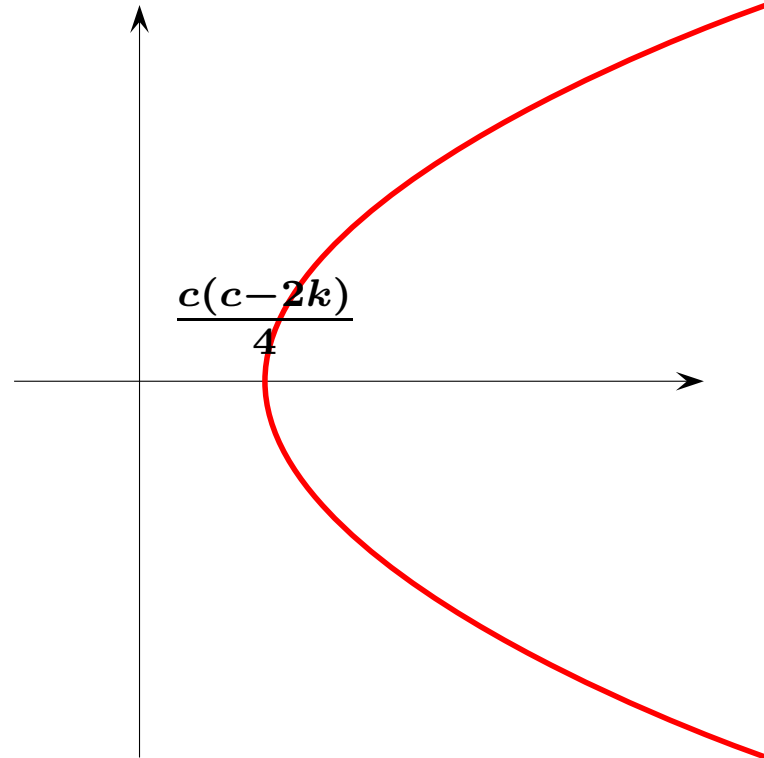
$$\sigma(-\mathfrak{A}) = \left[\frac{c^2}{4}, \infty\right)$$

and

$$\sigma(-\mathfrak{A} - b) = \left\{\frac{c(c-k)}{2} + \xi^2 + ik\xi; \xi \in \mathbb{R}\right\}.$$



$$-\mathfrak{A}$$



$$-\mathfrak{A} - k \frac{d}{dx}$$

Now we take a different point of view.

We fix an operator $\mathfrak{A} = \frac{d^2}{dx^2} - c \frac{d}{dx}$ but we change a reference measure. For $\theta \in [0, 1]$, define

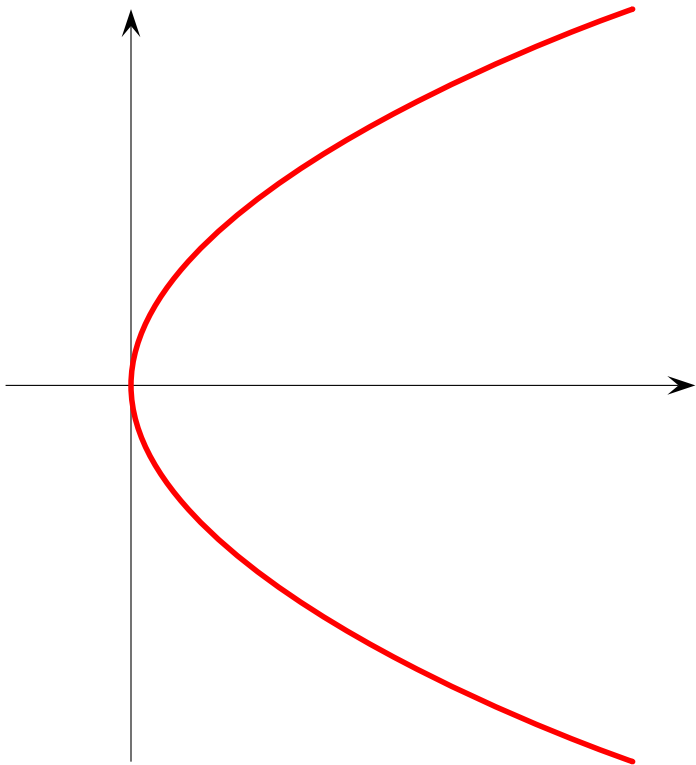
$$\nu_\theta(dx) = (1 - \theta)dx + \theta e^{-cx} dx$$

ν_θ is an invariant measure for \mathfrak{A} . $\nu_0(dx) = dx$, $\nu_1(dx) = e^{-cx} dx$.

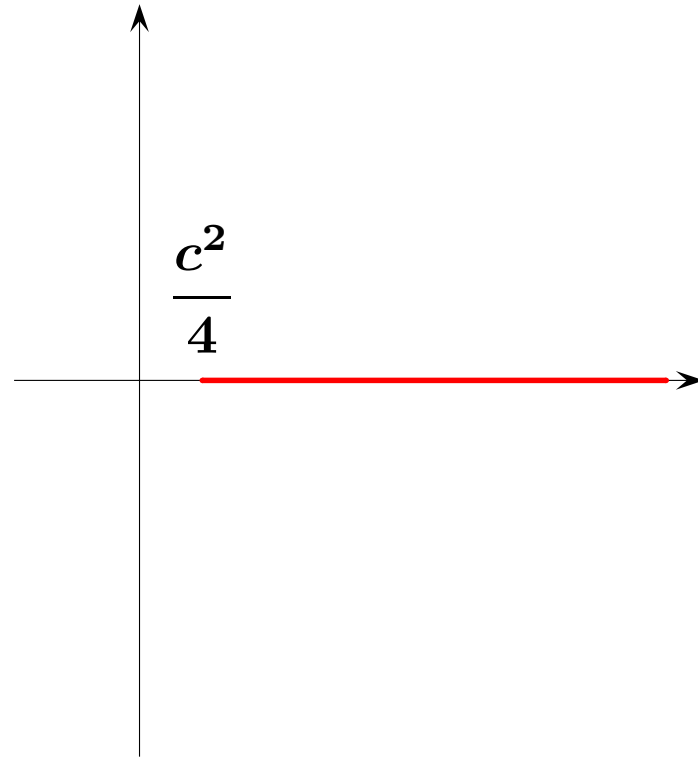
The computation above implies

$$\sigma(-\mathfrak{A}) = \{\xi^2 - ic\xi; \xi \in \mathbb{R}\} \quad \text{in } L^2(\nu_0)$$

$$\sigma(-\mathfrak{A}) = \left[\frac{c^2}{4}, \infty\right) \quad \text{in } L^2(\nu_1).$$



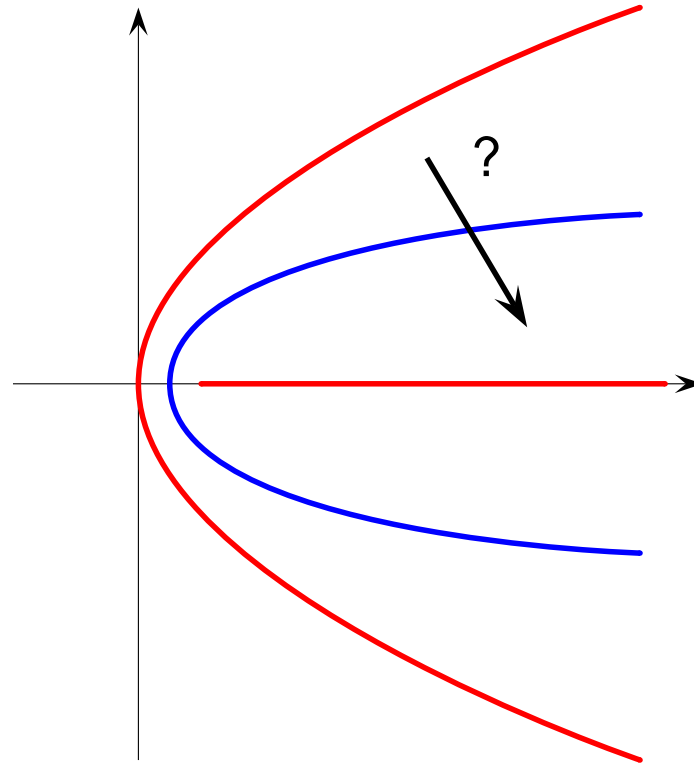
w.r.t. $\nu_0 = dx$



w.r.t. $\nu_1 = e^{-cx} dx$

What happens if we take the measure ν_θ ?

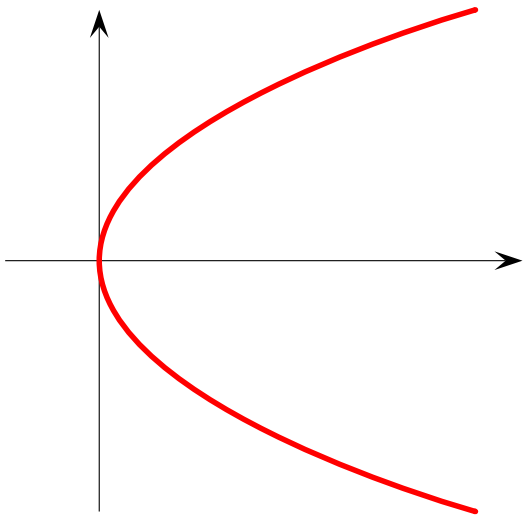
Does the spectrum change continuously?



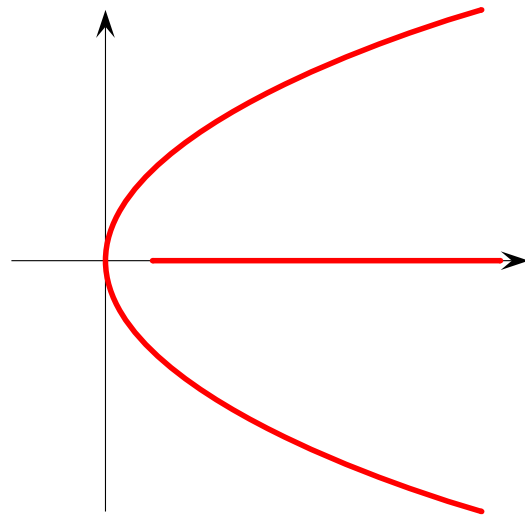
w.r.t. ν_θ

Theorem 8. For $\theta \in (0, 1)$, $\sigma(-\mathfrak{A})$ in $L^2(\nu_\theta)$ is

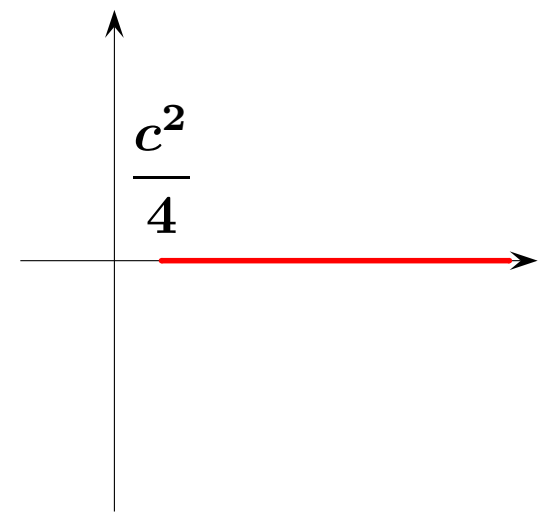
$$\{\xi^2 - ik\xi; \xi \in \mathbb{R}\} \cup \left[\frac{c^2}{4}, \infty\right).$$



w.r.t. $\nu_0 = dx$



w.r.t. ν_θ



w.r.t. $\nu_1 = e^{-cx} dx$

5. Perturbation by rotation

Laplacian on \mathbb{R}^2

Let \mathfrak{A} be

$$(5) \quad \mathfrak{A} = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + k \left(x \frac{\partial}{\partial y} - y \frac{\partial}{\partial x} \right) \quad \text{on } L^2(\mathbb{R}^2, dx dy).$$

The spectrum of $-\frac{\partial^2}{\partial x^2} - \frac{\partial^2}{\partial y^2}$ is $[0, \infty)$.

For the spectrum of \mathfrak{A} , we recall the Bessel functions:

$$J_\nu(x) = \left(\frac{x}{2} \right)^\nu \sum_{l=0}^{\infty} \frac{(ix/2)^{2l}}{l! \Gamma(\nu + l + 1)}, \quad \Re \nu > 0$$

which satisfies the following differential equation

$$I'' + \frac{1}{x} I' + \left(1 - \frac{\nu^2}{x^2} \right) I = 0$$

Since our space is \mathbb{R}^2 , we only need the case that ν is a non-negative integer. We use the polar coordinate:

$$\begin{cases} x = r \cos \theta, \\ y = r \sin \theta, \end{cases} \quad r \geq 0, \theta \in [0, 2\pi)$$

Using this, \mathfrak{A} can be written as

$$\mathfrak{A} = \frac{\partial^2}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r} + \frac{1}{r^2} \frac{\partial^2}{\partial \theta^2} + k \frac{\partial}{\partial \theta}.$$

If $F = f(r)e^{in\theta}$, then

$$\mathfrak{A}F = \left(\frac{\partial^2}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r} - n^2 \frac{1}{r^2} \right) f(r)e^{in\theta} + ikn f(r)e^{in\theta}.$$

Further

$$\left(\frac{\partial^2}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r} - n^2 \frac{1}{r^2} \right) J_n(\lambda r) = -\lambda^2 J_n(\lambda r).$$

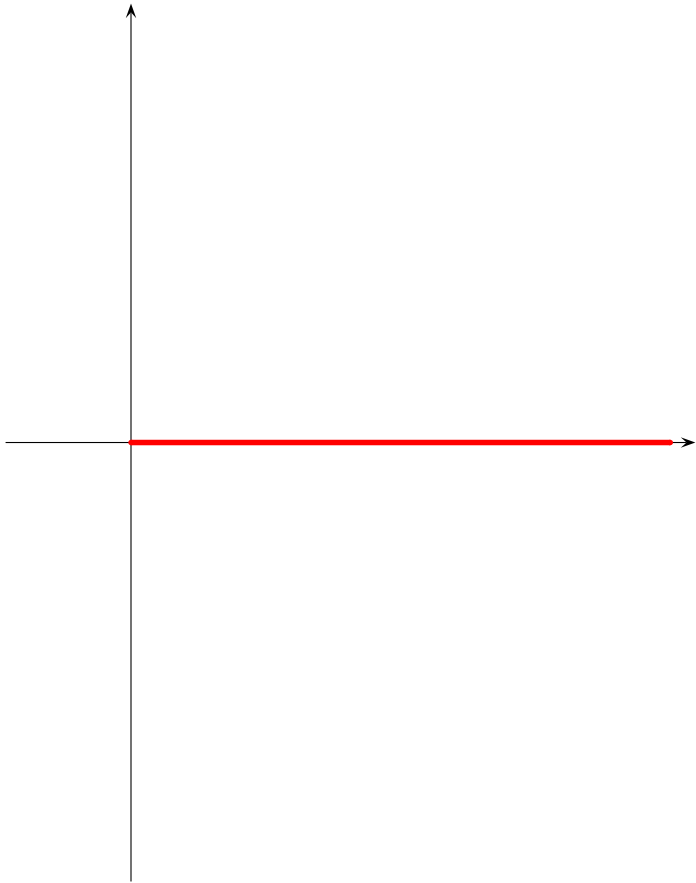
The spectral decomposition is given by

$$f(r, \theta) = \sum_{n \in \mathbb{Z}} \int \left\{ \int \left(\frac{1}{2\pi} \int f(\rho, \phi) e^{-in\phi} d\phi \right) J_{|n|}(\lambda\rho) \rho d\rho \right\} e^{in\theta} J_{|n|}(\lambda r) \lambda d\lambda.$$

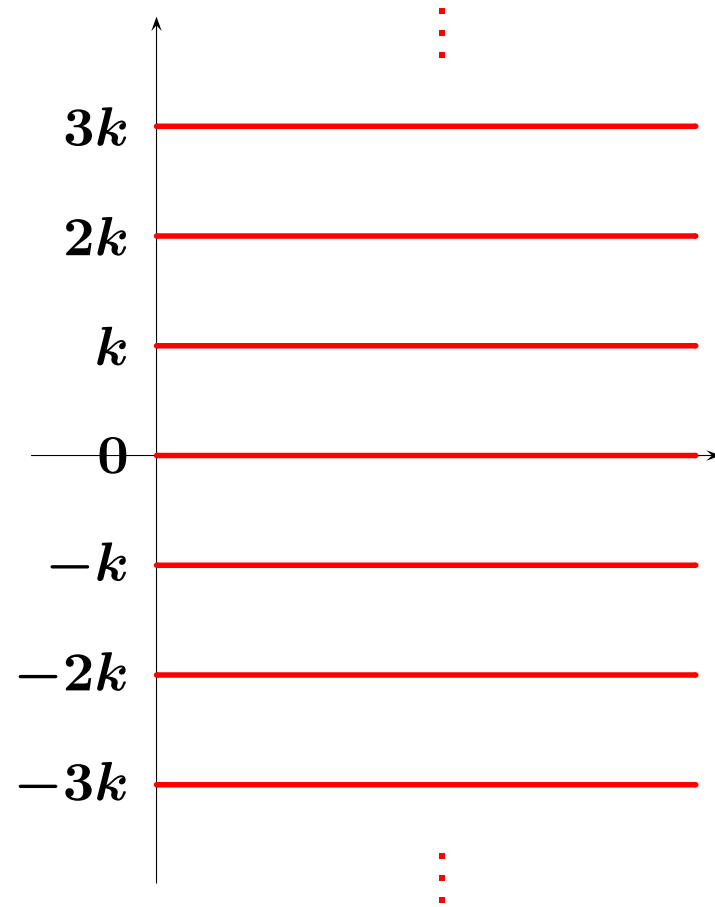
Theorem 9. The spectrum of $-\mathfrak{A}$ is

$$(6) \quad \{\lambda^2 - ikn; \lambda \geq 0, n \in \mathbb{Z}\}$$

and the corresponding eigenfunction to $\lambda^2 - ikn$ is $J_{|n|}(\lambda r) e^{in\theta}$.



the spectrum of $-\frac{\partial^2}{\partial x^2} - \frac{\partial^2}{\partial y^2}$



the spectrum of \mathfrak{A}

Ornstein Uhlenbeck operator on \mathbb{R}^2

Let L_α be

$$(7) \quad L_\alpha = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} - x \frac{\partial}{\partial x} - y \frac{\partial}{\partial y} + \alpha \left(x \frac{\partial}{\partial y} - y \frac{\partial}{\partial x} \right)$$

acting on $L^2(\mathbb{R}^2, \frac{1}{2\pi} e^{-(x^2+y^2)/2} dx dy)$.

The spectrum of Ornstein-Uhlenbeck operator L_0 is $\{0, -1, -2, \dots\}$. In fact, define Hermite polynomials by

$$H_n(x) = (-1)^n e^{x^2/2} \frac{d^n}{dx^n} e^{-x^2/2}.$$

Then

$$L_0 H_k(x) H_{n-k}(y) = -n H_k(x) H_{n-k}(y).$$

To get the spectram of L_α , we need the complex Hermite polynomials defined by

$$(8) \quad H_{p,q}(z, \bar{z}) = (-1)^{p+q} e^{\frac{z\bar{z}}{2}} \left(\frac{\partial}{\partial \bar{z}} \right)^p \left(\frac{\partial}{\partial z} \right)^q e^{-\frac{z\bar{z}}{2}}.$$

Here, we regard \mathbb{R}^2 as \mathbb{C} with $z = x + iy$. We denote

$$\frac{\partial}{\partial z} = \frac{1}{2} \left(\frac{\partial}{\partial x} - i \frac{\partial}{\partial y} \right), \quad \frac{\partial}{\partial \bar{z}} = \frac{1}{2} \left(\frac{\partial}{\partial x} + i \frac{\partial}{\partial y} \right).$$

In the sequel, we write

$$\partial = \frac{\partial}{\partial z}, \quad \bar{\partial} = \frac{\partial}{\partial \bar{z}}$$

for short. We have

$$(9) \quad \partial^* = -\bar{\partial} + \frac{z}{2}, \quad \bar{\partial}^* = -\partial + \frac{\bar{z}}{2}$$

Proposition 10. The following identities hold:

$$\partial H_{p,q} = \frac{p}{2} H_{p-1,q}, \quad \bar{\partial} H_{p,q} = \frac{q}{2} H_{p,q-1},$$

$$\partial^* H_{p,q} = H_{p+1,q}, \quad \bar{\partial}^* H_{p,q} = H_{p,q+1},$$

$$(2\partial\bar{\partial} - z\partial)H_{p,q} = -pH_{p,q}$$

$$(2\partial\bar{\partial} - \bar{z}\bar{\partial})H_{p,q} = -qH_{p,q}$$

$$(z\partial - \bar{z}\bar{\partial})H_{p,q} = (p - q)H_{p,q}$$

We can write

$$L_\alpha = (2\partial\bar{\partial} - z\partial) + (2\partial\bar{\partial} - \bar{z}\bar{\partial}) + \alpha i(z\partial - \bar{z}\bar{\partial}).$$

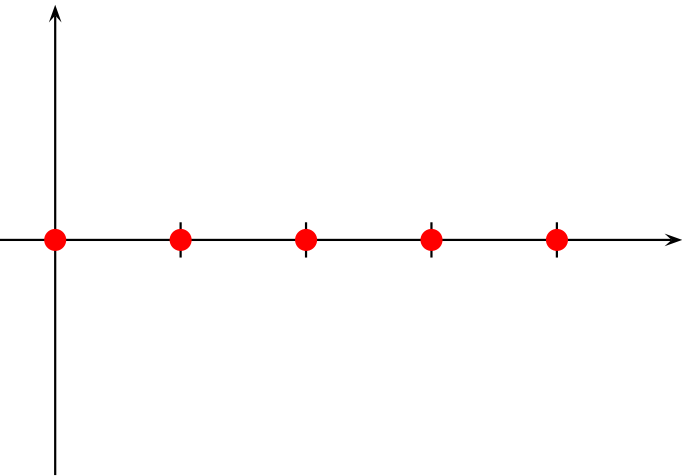
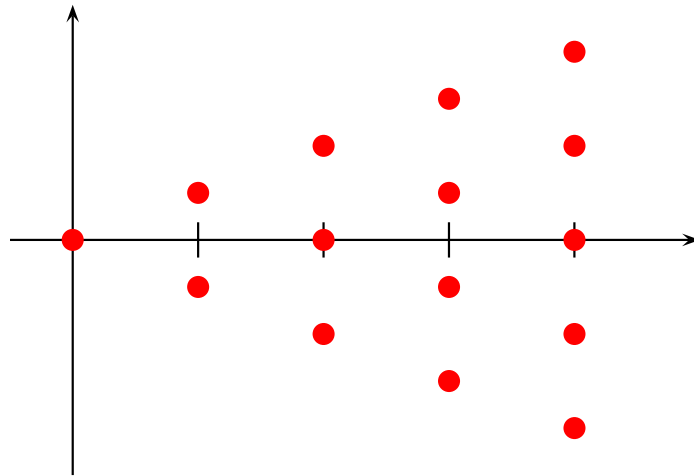
Hence

$$L_\alpha H_{p,q} = (-p - q + (p - q)\alpha i)H_{p,q}.$$

Theorem 11. The spectrum of $-L_\alpha$ is

$$(10) \quad \{(p + q) - (p - q)\alpha i\}_{p,q=0}^{\infty}$$

and corresponding eigenfunctions are $H_{p,q}$ respectively.

the spectrum of $-L_0$ the spectrum of $-L_\alpha$

Connection to the Laguerre polynomials

The eigenfunction $H_{n,n}$ for the eigenvalue $2n$, ($n \in \mathbb{Z}_+$) is rotation invariant since $\left(x \frac{\partial}{\partial y} - y \frac{\partial}{\partial x}\right) H_{n,n} = 0$. So $H_{n,n}$ is a function of $r = |z|$ and

$$\left(\frac{d^2}{dr^2} + \frac{1}{r} \frac{d}{dr} - r \frac{d}{dr}\right) H_{n,n} = -2n H_{n,n}.$$

Now, by the change of variable $r = \sqrt{2u}$, we have

$$\frac{d^2}{dr^2} + \frac{1}{r} \frac{d}{dr} - r \frac{d}{dr} = 2u \frac{d^2}{du^2} + 2(1-u) \frac{d}{du}.$$

$F(u) = H_{n,n}(r)$ satisfies

$$2u \frac{d^2}{du^2} F + 2(1-u) \frac{d}{du} F + nF = 0.$$

The Laguerre polynomial satisfies this differential equation. Here the Laguerre polynomial polynomial is defined by

$$(11) \quad L_n = \frac{e^x}{n!} \frac{d^n}{dx^n} (e^{-x} x^n)$$

Now we have

Theorem 12. Complex Hermite polynomials $H_{n,n}$ are expressed as following;

$$(12) \quad H_{n,n}(z, \bar{z}) = \frac{(-1)^n n!}{2^n} L_n \left(\frac{|z|^2}{2} \right)$$

where c is a constant.

Thanks a lot!