

On the dichotomy in the heat kernel two sided estimates

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1 Preliminaries

Let (M, d) be a locally compact separable metric space and μ be a Radon measure on M with full support. Assume that there exists a heat kernel $\{p_t\}_{t>0}$ on M :

Definition. A family $\{p_t\}_{t>0}$ of Borel functions $p_t(x, y)$ on $M \times M$ is called a *heat kernel* if the following conditions are satisfied, for all $x, y \in M$ and all $s, t > 0$:

- (i) Positivity: $p_t(x, y) \geq 0$.
- (ii) The total mass inequality

$$\int_M p_t(x, y) d\mu(y) \leq 1. \quad (1.1)$$

(iii) Symmetry: $p_t(x, y) = p_t(y, x)$.

(iv) Semigroup property:

$$p_{s+t}(x, y) = \int_M p_s(x, z)p_t(z, y)d\mu(z). \quad (1.2)$$

(v) Approximation of identity: for any $u \in L^2(M, \mu)$

$$\int_M p_t(x, y)u(y)d\mu(y) \xrightarrow{L^2} u(x) \quad \text{as } t \rightarrow 0+. \quad (1.3)$$

Any heat kernel gives rise to the *heat semigroup* $\{P_t\}_{t>0}$ where P_t is the operator in L^2 defined by

$$P_t u(x) = \int_M p_t(x, y)u(y)d\mu(y). \quad (1.4)$$

The semigroup identity (1.2) implies that $P_t P_s = P_{t+s}$, that is, the family $\{P_t\}_{t>0}$ is a *semigroup*. It follows from (1.3) that

$$s\text{-}\lim_{t \rightarrow 0} P_t = I,$$

where I is the identity operator in L^2 and $s\text{-}\lim$ stands for *strong* limit. Hence, $\{P_t\}_{t>0}$ is a strongly continuous, self-adjoint, contraction semigroup in L^2 .

Given the semigroup $\{P_t\}_{t>0}$, define the *infinitesimal generator* \mathcal{L} of the semigroup by

$$\mathcal{L}u := \lim_{t \rightarrow 0} \frac{u - P_t u}{t}, \quad (1.5)$$

where the limit is understood in the L^2 -norm. The *domain* $\text{dom}(\mathcal{L})$ of the generator \mathcal{L} is the space of functions $u \in L^2$ for which the limit in (1.5) exists. By the Hille–Yosida theorem, $\text{dom}(\mathcal{L})$ is dense in L^2 . Furthermore, \mathcal{L} is a self-adjoint, positive definite operator, which immediately follows from the fact that the semigroup $\{P_t\}$ is self-adjoint and contractive. Moreover, we have

$$P_t = \exp(-t\mathcal{L}), \quad (1.6)$$

where the right hand side is understood in the sense of spectral theory.

Let a heat kernel p_t on (M, d, μ) satisfy the following two-sided estimate for all $x, y \in M$ and all $t \in (0, \infty)$:

$$\frac{c_1}{t^{\alpha/\beta}} \Phi\left(c_2 \frac{d(x, y)}{t^{1/\beta}}\right) \leq p_t(x, y) \leq \frac{c_3}{t^{\alpha/\beta}} \Phi\left(c_4 \frac{d(x, y)}{t^{1/\beta}}\right) \quad (1.7)$$

where α, β are positive constants and Φ is a non-negative monotone decreasing function on $[0, \infty)$. There are two very important classes of heat kernels that satisfy (1.7). One is the heat kernel of diffusions on various fractals, where the function Φ is of the form

$$\Phi(s) = \exp(-s^\gamma), \quad (1.8)$$

for some $\gamma > 0$ (see [2] and the references therein). The Gauss-Weierstrass heat kernel on \mathbb{R}^d is included in this class, in which case $\alpha = d, \beta = \gamma = 2$. The other is the heat kernel of stable-like processes, where the function Φ is of the form

$$\Phi(s) = (1 + s)^{-\alpha-\beta}, \quad (1.9)$$

(see [5] and the references therein). The heat kernel of the symmetric β -stable process on \mathbb{R}^d is included in this class, in which case $\alpha = d, 0 < \beta < 2$.

The nature of the parameters α and β is important. The parameter α turns out to be the Hausdorff dimension of M . The parameter β is called the walk dimension of the heat kernel p_t . This terminology is from the following observation: if the heat kernel p_t is the transition density of a Markov process X_t on M , then under mild assumptions about Φ , (1.7) implies that the average time t needed for the process X_t to move away to a distance r from the origin is of the order r^β (see, for example, [2, Lemma 3.9]).

As mentioned above, there are important classes of heat kernels that satisfy (1.7). One can then ask the following natural question:

Is there a heat kernel p_t satisfying (1.7), where Φ is different from (1.8) and (1.9)?

The main purpose of this paper is to answer this question. In Theorem 3.4, we will show under mild assumptions that the shape of Φ for any heat kernel p_t satisfying (1.7) is either (1.8) or (1.9). Our approach is analytic, which does not depend on the existence of the process X_t .

2 Dirichlet form associated with a heat kernel

Let (M, d, μ) be a metric measure space with a heat kernel $\{p_t\}_{t>0}$, and let $\{P_t\}_{t>0}$ be the heat semigroup defined by (1.4). For any $t > 0$, we define a quadratic form \mathcal{E}_t on L^2 by

$$\mathcal{E}_t[u] := \left(\frac{u - P_t u}{t}, u \right), \quad (2.10)$$

where (\cdot, \cdot) is the inner product in L^2 . An easy computation shows that \mathcal{E}_t can be equivalently defined by

$$\mathcal{E}_t[u] = \frac{1}{2t} \int_M \int_M |u(x) - u(y)|^2 p_t(x, y) d\mu(y) d\mu(x). \quad (2.11)$$

In terms of the spectral resolution $\{E_\lambda\}$ of the generator \mathcal{L} , \mathcal{E}_t can be expressed as follows

$$\mathcal{E}_t[u] = \int_0^\infty \frac{1 - e^{-t\lambda}}{t} d\|E_\lambda u\|_2^2,$$

which implies that $\mathcal{E}_t[u]$ is decreasing in t (indeed, this is an elementary exercise to show that the function $t \mapsto \frac{1 - e^{-t\lambda}}{t}$ is decreasing).

Let us define a quadratic form \mathcal{E} by

$$\mathcal{E}[u] := \lim_{t \rightarrow 0^+} \mathcal{E}_t[u] = \int_0^\infty \lambda d\|E_\lambda u\|_2^2 \quad (2.12)$$

(where the limit may be $+\infty$ since $\mathcal{E}[u] \geq \mathcal{E}_t[u]$) and its domain $\mathcal{D}(\mathcal{E})$ by

$$\mathcal{D}(\mathcal{E}) := \{u \in L^2 : \mathcal{E}[u] < \infty\}.$$

It is clear from (2.11) and (2.12) that \mathcal{E}_t and \mathcal{E} are positive definite.

It is easy to see from (2.12) that $\mathcal{D}(\mathcal{E}) = \text{dom}(\mathcal{L}^{1/2})$. The domain $\mathcal{D}(\mathcal{E})$ is dense in L^2 because $\mathcal{D}(\mathcal{E})$ contains $\text{dom}(\mathcal{L})$. Indeed, if $u \in \text{dom}(\mathcal{L})$ then using (1.5) and (2.10), we obtain

$$\mathcal{E}[u] = \lim_{t \rightarrow 0} \mathcal{E}_t[u] = (\mathcal{L}u, u) < \infty. \quad (2.13)$$

The quadratic form $\mathcal{E}[u]$ extends to a bilinear form $\mathcal{E}(u, v)$ by the polarization identity

$$\mathcal{E}(u, v) = \frac{1}{2}(\mathcal{E}[u+v] - \mathcal{E}[u-v]).$$

It follows from (2.13) that $\mathcal{E}(u, v) = (\mathcal{L}u, v)$ for all $u, v \in \text{dom}(\mathcal{L})$. The space $\mathcal{D}(\mathcal{E})$ is naturally endowed with the inner product

$$[u, v] := (u, v) + \mathcal{E}(u, v). \quad (2.14)$$

It is possible to show that the form \mathcal{E} is *closed*, that is the space $\mathcal{D}(\mathcal{E})$ is *Hilbert*.

It is easy to see from (1.4) and the definition of a heat kernel that the semigroup $\{P_t\}$ is *Markovian*, that is $0 \leq u \leq 1$ implies $0 \leq P_t u \leq 1$. This implies that the form \mathcal{E} satisfies the Markov property, that is $u \in \mathcal{D}(\mathcal{E})$ implies $v := \min(u_+, 1) \in \mathcal{D}(\mathcal{E})$ and $\mathcal{E}[v] \leq \mathcal{E}[u]$. Hence, \mathcal{E} is a *Dirichlet form*.

We say that \mathcal{E} is *local* if $\text{supp } u$ and $\text{supp } v$ are disjoint compact sets for $u, v \in \mathcal{D}(\mathcal{E})$, then $\mathcal{E}(u, v) = 0$. \mathcal{E} is called *stochastically complete* if $P_t 1 = 1$ for all $t > 0$, that is, the equality holds in (1.1).

A Dirichlet form \mathcal{E} is said to be *regular* if there exists a subspace $\mathcal{C} \subset \mathcal{D}(\mathcal{E}) \cap C_0(M)$ such that \mathcal{C} is dense in $\mathcal{D}(\mathcal{E})$ with $[\cdot]$ -norm and dense in $C_0(M)$ with uniform norm. (Here $C_0(M)$ is the space of continuous compactly supported functions on M .) When \mathcal{E} is *regular*, there is a corresponding Markov process X_t which is furthermore a Hunt process. As we mentioned in the first section, we do *not* assume \mathcal{E} to be regular throughout this paper.

3 Main result

Fix two positive parameters α and β and a monotone decreasing function $\Phi : [0, +\infty) \rightarrow [0, +\infty)$ such that $\Phi(c) > 0$ for some $c > 0$.

Lemma 3.1 *Assume that $\{p_t\}$ is a heat kernel on (M, d, μ) such that, for all $x, y \in M$ and $t > 0$,*

$$p_t(x, y) \leq \frac{C}{t^{\alpha/\beta}} \Phi\left(\frac{d(x, y)}{t^{1/\beta}}\right), \quad (3.1)$$

for some $C > 0$. Then either the associated Dirichlet form \mathcal{E} is local or

$$\Phi(s) \geq \frac{c}{(1+s)^{\alpha+\beta}} \quad (3.2)$$

for some $c > 0$.

Proof. Consider the form \mathcal{E}_t which is given by

$$\mathcal{E}_t(u, v) = \frac{1}{2t} \int_M \int_M (u(x) - u(y))(v(x) - v(y)) p_t(x, y) d\mu(y) d\mu(x). \quad (3.3)$$

Let $u, v \in L^1(M, \mu)$ be two non-negative functions with disjoint supports $A = \text{supp } u$ and $B = \text{supp } v$, and set

$$r = d(A, B) > 0. \quad (3.4)$$

The integrand in (3.3) vanishes if either both x, y are outside A or both x, y are outside B . Hence, we can restrict the integration to the domain where one of the variables x, y is in A and the other is in B . Hence, we obtain, using the symmetry of the heat kernel,

$$\begin{aligned} \mathcal{E}_t(u, v) &= -\frac{1}{2t} \int_A \int_B u(x)v(y) p_t(x, y) d\mu(y) d\mu(x) \\ &\quad -\frac{1}{2t} \int_B \int_A u(y)v(x) p_t(x, y) d\mu(y) d\mu(x) \\ &= -\frac{1}{t} \int_A \int_B u(x)v(y) p_t(x, y) d\mu(y) d\mu(x). \end{aligned} \quad (3.5)$$

If $x \in A$ and $y \in B$ then $d(x, y) \geq r$ and

$$p_t(x, y) \leq \frac{C}{t^{\alpha/\beta}} \Phi\left(\frac{r}{t^{1/\beta}}\right).$$

Therefore, (3.5) implies

$$|\mathcal{E}_t(u, v)| \leq \frac{C}{t^{1+\alpha/\beta}} \Phi\left(\frac{r}{t^{1/\beta}}\right) \|u\|_{L^1} \|v\|_{L^1}. \quad (3.6)$$

If (3.2) fails then there exists a sequence $\{s_k\} \rightarrow \infty$ such that

$$s_k^{\alpha+\beta} \Phi(s_k) \rightarrow 0 \quad \text{as } k \rightarrow \infty.$$

Define a sequence $\{t_k\}$ from the condition

$$s_k = \frac{r}{t_k^{1/\beta}}.$$

Then

$$s_k^{\alpha+\beta} \Phi(s_k) = \frac{r^{\alpha+\beta}}{t_k^{1+\alpha/\beta}} \Phi\left(\frac{r}{t_k^{1/\beta}}\right) \rightarrow 0 \quad \text{as } k \rightarrow \infty,$$

and (3.6) implies that

$$\mathcal{E}_{t_k}(u, v) \rightarrow 0 \quad \text{as } k \rightarrow \infty. \quad (3.7)$$

Therefore, if $\text{supp } u$ and $\text{supp } v$ are disjoint compact sets for $u, v \in \mathcal{D}(\mathcal{E})$, then we can take $r > 0$ as in (3.4), so by (3.7), $\mathcal{E}(u, v) = \lim_{k \rightarrow \infty} \mathcal{E}_{t_k}(u, v) = 0$. Hence the locality of \mathcal{E} follows. ■

Lemma 3.2 Assume that $\{p_t\}$ is a heat kernel on (M, d, μ) such that, for all $x, y \in M$ and $t > 0$,

$$p_t(x, y) \geq \frac{c}{t^{\alpha/\beta}} \Phi\left(\frac{d(x, y)}{t^{1/\beta}}\right), \quad (3.8)$$

for some $c > 0$. Then

$$\Phi(s) \leq \frac{C}{(1+s)^{\alpha+\beta}} \quad (3.9)$$

for some $C > 0$.

Proof. Let u be a non-constant function from $L^2(M, \mu)$. Choose a ball $Q \subset M$ where u is non-constant and let $a > b$ be two real values such that the sets

$$A = \{x \in Q : u(x) \geq a\} \text{ and } B = \{x \in Q : u(x) \leq b\}$$

have positive measures. If the diameter of Q is D then, by (3.8), we have for all $x, y \in Q$

$$p_t(x, y) \geq \frac{c}{t^{\alpha/\beta}} \Phi\left(\frac{D}{t^{1/\beta}}\right)$$

whence by (3.3)

$$\begin{aligned} \mathcal{E}(u, u) &\geq \mathcal{E}_t(u, u) \geq \frac{1}{2t} \int_A \int_B (u(x) - u(y))^2 p_t(x, y) d\mu(y) d\mu(x) \\ &\geq (a - b)^2 \mu(A) \mu(B) \frac{c}{2t^{1+\alpha/\beta}} \Phi\left(\frac{D}{t^{1/\beta}}\right) \\ &= \frac{c'}{t^{1+\alpha/\beta}} \Phi\left(\frac{D}{t^{1/\beta}}\right) \end{aligned} \quad (3.10)$$

where $c' > 0$. If (3.9) fails then there exists a sequence $\{s_k\} \rightarrow \infty$ such that

$$s_k^{\alpha+\beta} \Phi(s_k) \rightarrow \infty \text{ as } k \rightarrow \infty.$$

Define a sequence $\{t_k\}$ from the condition

$$s_k = \frac{D}{t_k^{1/\beta}}.$$

Then

$$\frac{1}{t_k^{1+\alpha/\beta}} \Phi\left(\frac{D}{t_k^{1/\beta}}\right) = D^{-(\alpha+\beta)} s_k^{\alpha+\beta} \Phi(s_k) \rightarrow \infty \text{ as } k \rightarrow \infty,$$

and (3.10) yields $\mathcal{E}(u, u) = \infty$. Therefore, the domain of the form \mathcal{E} contains only constants. Note that $L^2(M, \mu)$ does not consist of only constants. (Indeed, since μ is a Radon measure on M with full support, it is enough to check that M consists of more than one point. By (3.8), $p_t(x, x) \rightarrow \infty$ as $t \rightarrow 0$, so if $M = \{x\}$, this contradicts (1.1).) Thus, the fact that the form \mathcal{E} contains only constants contradicts the fact that this domain is dense in $L^2(M, \mu)$. ■

We say that (M, d) satisfies the chain condition if there exists a (large) constant C such that for any two points $x, y \in M$ and for any positive integer n , there exists a sequence $\{x_i\}_{i=0}^n$ of points in M such that $x_0 = x$, $x_n = y$, and

$$d(x_i, x_{i+1}) \leq C \frac{d(x, y)}{n}, \quad \text{for all } i = 0, 1, \dots, n-1.$$

In the following, we denote $\Phi(s) \simeq f(s)$ if there exist constants $c_1, c_2 > 0$ such that $c_1 f(s) \leq \Phi(s) \leq c_2 f(s)$ for all $s > 0$. Similarly, we denote $\Phi(s) \asymp f(Cs)g(cs)$ if there exist constants $c_1, \dots, c_4 > 0$ such that $f(c_1 s)g(c_2 s) \leq \Phi(s) \leq f(c_3 s)g(c_4 s)$ for all $s > 0$.

Corollary 3.3 *If the following estimate holds for all $x, y \in M$ and $t > 0$,*

$$p_t(x, y) \asymp \frac{C}{t^{\alpha/\beta}} \Phi \left(c \frac{d(x, y)}{t^{1/\beta}} \right) \quad (3.11)$$

then either the Dirichlet form \mathcal{E} is local or

$$\Phi(s) \simeq \frac{1}{(1+s)^{\alpha+\beta}}. \quad (3.12)$$

Proof. Indeed, if \mathcal{E} is non-local then Φ satisfies (3.2) and (3.9), whence the claim follows. ■

Theorem 3.4 *Let the metric space (M, d) satisfy the chain condition, the heat kernel be stochastically complete, and (3.11) hold with some $\alpha, \beta > 0$ and Φ . Then $\beta \leq \alpha + 1$,*

$$\mu(B(x, r)) \simeq r^\alpha, \quad (3.13)$$

and the following dichotomy holds:

- *either the Dirichlet form \mathcal{E} is local, $\beta \geq 2$, and $\Phi(s) \asymp C \exp\left(-cs^{\frac{\beta}{\beta-1}}\right)$.*
- *or the Dirichlet form \mathcal{E} is non-local and $\Phi(s) \simeq (1+s)^{-(\alpha+\beta)}$.*

Proof. By Lemma 3.2, we have the upper bound

$$\Phi(s) \leq \frac{C}{(1+s)^{\alpha+\beta}}. \quad (3.14)$$

In particular, this implies

$$\int_0^\infty s^{\alpha-1} \Phi(s) ds < \infty. \quad (3.15)$$

By [9, Theorem 3.2] (see also [6]), the estimate (3.11) with a function Φ satisfying (3.15) and the stochastic completeness imply (3.13). By [9, Corollary 3.3], we have $\text{diam}(M) = \infty$. Furthermore, as it follows from the proof of [9, Theorem 3.2], for any $\varepsilon > 0$ there is $\delta > 0$ such that

$$\int_{B(x, r)^c} p_t(x, y) d\mu(y) \leq \varepsilon, \quad (3.16)$$

provided $t \leq (\delta r)^\beta$.

Also, (3.11) with (3.15) and the chain condition imply that $\beta \leq \alpha + 1$ (see [9, Theorem 4.8(ii)] and [6])

If the form \mathcal{E} is non-local, then by Lemma 3.1 and Lemma 3.2, Φ satisfies (3.12), which finishes the proof in this case. Assume now that \mathcal{E} is local. In this case, we will show that (3.16) implies that $\beta \geq 1$ and, for all $t, r > 0$ and $x \in M$,

$$\int_{B(x,r)^c} p_t(x,y) d\mu(y) \leq \begin{cases} C \exp\left(-c \left(\frac{r^\beta}{t}\right)^{\frac{1}{\beta-1}}\right), & \text{if } \beta > 1, \\ C \exp\left(-c \left(\frac{r}{t}\right)\right), & \text{if } \beta = 1. \end{cases} \quad (3.17)$$

Indeed, for each $\beta > 0$, using (4.21) and (4.22) in [8], letting $k \rightarrow \infty$ and then replacing $2r$ by r , we have

$$\int_{B(x,r)^c} p_t(x,y) d\mu(y) \leq C \exp\left(\lambda t - c_1 \lambda^{1/\beta} r\right), \quad \text{for all } \lambda \geq c_2 r^{-\beta}. \quad (3.18)$$

(Note that the arguments in [8] do not require the regularity of the Dirichlet form.) When $\beta < 1$, take $\lambda = c_3(r/t)^{\beta/(\beta-1)}$, where $c_3 - c_1 c_3^{1/\beta} = -1$. Then, $\lambda \geq c_2 r^{-\beta}$ is equivalent to $t \geq c_4 r^\beta$ for some $c_4 > 0$, so we obtain

$$\int_{B(x,r)^c} p_t(x,y) d\mu(y) \leq C \exp\left(-\left(\frac{t}{r^\beta}\right)^{\frac{1}{1-\beta}}\right) \quad \text{for } t \geq c_4 r^\beta. \quad (3.19)$$

On the other hand, by the lower bound of (3.11), for $t = Mr^\beta$ and $y \in B(x, 2r)$, we have

$$p_t(x,y) \geq \frac{C}{t^{\alpha/\beta}} \Phi\left(c \frac{2r}{t^{1/\beta}}\right) = \frac{C}{M^{\alpha/\beta} r^\alpha} \Phi\left(\frac{2c}{M^{1/\beta}}\right) \geq \frac{C'}{M^{\alpha/\beta} r^\alpha} \quad (3.20)$$

when M is large enough, since Φ is monotone decreasing and $\Phi(a) > 0$ for some $a > 0$. Integrating (3.20) over $y \in B(x, 2r) \setminus B(x, r)$ and using (3.13), we have that the left hand side of (3.19) is greater than or equal to $C'' M^{-\alpha/\beta}$. This is a contradiction when M is very large, because the right hand side of (3.19) is $C \exp(-M^{1/(1-\beta)})$ for $t = Mr^\beta$. So, we obtain $\beta \geq 1$. Now, applying (3.18) with $\lambda = c(r/t)^{\beta/(\beta-1)}$ when $\beta > 1$ and with $\lambda = ct^{-1}$ when $\beta = 1$, we obtain (3.17) for $t \leq c'r^\beta$. (3.17) is always true for $t \geq c'r^\beta$ by adjusting C , so the proof of (3.17) is completed.

Now, $x, y \in M$, $t > 0$, and for $r = \frac{1}{2}d(x, y)$,

$$\begin{aligned} p_t(x,y) &= \int_M p_{t/2}(x,z) p_{t/2}(z,y) d\mu(z) \\ &\leq \int_{B(x,r)^c \cup B(y,r)^c} p_{t/2}(x,z) p_{t/2}(z,y) d\mu(z) \\ &\leq \sup_{z \in M} p_{t/2}(z,y) \int_{B(x,r)^c} p_{t/2}(x,z) d\mu(z) + \sup_{z \in M} p_{t/2}(x,z) \int_{B(y,r)^c} p_{t/2}(y,z) d\mu(z). \end{aligned}$$

Since by (3.11) $p_t(x,y) \leq Ct^{-\alpha/\beta}$ for all $x, y \in M$ and $t > 0$, combining this with (3.17) we obtain,

$$p_t(x,y) \leq \begin{cases} \frac{C}{t^{\alpha/\beta}} \exp\left(-c \left(\frac{d^\beta(x,y)}{t}\right)^{\frac{1}{\beta-1}}\right), & \text{if } \beta > 1, \\ \frac{C}{t^\alpha} \exp\left(-c \left(\frac{r}{t}\right)\right), & \text{if } \beta = 1. \end{cases} \quad (3.21)$$

Now, by [9, Theorem 4.8(i)], the estimates (3.11) and (3.21) imply $\beta \geq 2$. (Note that the arguments in [9] do not require the regularity of the Dirichlet form.)

On the other hand, when $\beta \geq 2$ (in fact $\beta > 1$ is enough), the standard chain argument (see [9, Corollary 3.5]) shows that the lower bound in (3.11) implies the lower bounds

$$p_t(x, y) \geq \frac{C}{t^{\alpha/\beta}} \exp\left(-c \left(\frac{d^\beta(x, y)}{t}\right)^{\frac{1}{\beta-1}}\right).$$

Combining these estimates with (3.11), we obtain

$$\Phi(s) \asymp C \exp\left(-cs^{\frac{\beta}{\beta-1}}\right),$$

with $\beta \geq 2$, which finishes the proof. ■

Remark. 1) This theorem excludes discrete cases. Indeed, for the discrete cases, (3.11) does not hold for very small t . For example, continuous time simple random walk on \mathbb{Z}^d satisfies (3.11) with $\alpha = d, \beta = 2$ and $\Phi(s) \asymp C \exp(-cs^2)$ for $d(x, y) \vee 1 \leq t$, but (3.11) does not hold for $t \ll 1$.

2) In this theorem, we assume (3.11) for all $x, y \in M$ and $t > 0$. But if the Dirichlet form \mathcal{E} is regular, then we can relax this part of the assumption and need only to assume (3.11) for μ -a.e. $x, y \in M$ and all $t > 0$. See [3, Theorem 2.1] and [4].

3) In the case of a local form, we obtain the relations between α and β

$$2 \leq \beta \leq \alpha + 1. \tag{3.22}$$

By [1], any couple of α, β in this range can be realized for the above heat kernel estimates. In the case of a non-local form, we have instead the range

$$0 < \beta \leq \alpha + 1.$$

Any couple in the range $0 < \beta < \alpha + 1$ can be realized. Indeed, if \mathcal{L} is the generator of diffusion with parameters α and β from the range (3.22) then $\mathcal{L}^\delta, \delta \in (0, 1)$, generated a Hunt process with the walk dimension $\beta' = \delta\beta$ and the same α , so that β' can take any value from $(0, \alpha + 1)$. We do not know whether $\beta = \alpha + 1$ can occur for non-local processes or not.

References

- [1] **Barlow M.T.**, Which values of the volume growth and escape time exponents are possible for graphs?, *Rev. Math. Iberoamericana*, **20** (2004) 1-31.
- [2] **Barlow M.T.**, Diffusions on fractals, *Lectures on Probability Theory and Statistics, Ecole d'Été de Probabilités de Saint-Flour XXV - 1995. Springer Lecture Notes Math.*, **1690** (1998) 1-121.
- [3] **Barlow M.T., Bass R.F., Chen Z.-Q., Kassmann M.**, Non-local Dirichlet forms and symmetric jump processes, preprint (2006)
- [4] **Barlow M.T., Grigor'yan A., Kumagai T.**, Heat kernel upper bounds for jump processes and the first exit time, preprint (2006)
- [5] **Chen Z.-Q., Kumagai T.**, Heat kernel estimates for stable-like processes on d -sets, *Stochastic Process Appl.*, **108** (2003) 27-62.

- [6] **Grigor'yan A.**, Heat kernels and function theory on metric measure spaces, *in*: “Heat kernels and analysis on manifolds, graphs, and metric spaces”, *Contemporary Mathematics*, **338** (2003) 143-172.
- [7] **Grigor'yan A.**, Heat kernel upper bounds on fractal spaces, preprint (2006)
- [8] **Grigor'yan A., Hu J.**, Upper estimates of transition densities for Dirichlet forms on metric spaces, in preparation(2007)
- [9] **Grigor'yan A., Hu J., Lau K.S.**, Heat kernels on metric-measure spaces and an application to semi-linear elliptic equations, *Trans. Amer. Math. Soc.*, **355** (2003) no.5, 2065-2095.