Witten Laplacian for an unbounded spin system^{*}

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We consider Witten Laplacians acting on differential forms on an unbounded lattice spin system. The configulation space is $\mathbb{R}^{\mathbb{Z}^d}$ and we are given a Gibbs measure on it. The Gibbs measure in this note is expressed formally as

$$\nu = Z^{-1} \exp\left\{-2\mathscr{J} \sum_{\substack{i,j \in \mathbb{Z}^d \\ i \sim j}} (x^i - x^j)^2 - 2\sum_{i \in \mathbb{Z}^d} U(x^i)\right\} \prod_{i \in \mathbb{Z}^d} dx^i.$$
(1)

Here \mathscr{I} is a positive constnat and U is an \mathbb{R} -valued function on \mathbb{R} and $i \sim j$ means that $|i-j|^2 = (i_1 - j_1)^2 + \cdots + (i_1 - j_1)^2 = 1$. Precise characterization is fomulated through Dobrushin-Lanford-Ruelle equation. We show that there is no harmonic *p*-forms ($p \geq 1$) on this space. To do this, we take a finite set $\Lambda \subset \mathbb{Z}^d$ and a boundary condition η and define a finite volume Gibbs measure on \mathbb{R}^{Λ} and then take limit.

For this model, the logarithmic Sobolev inequality [4, 3], the spectral gap and the vanishing theorem for 1-forms [1] under suitable condition, were already established. So we intend to generalize the vanishing theorem for general *p*-forms.

Finite volume Gibbs measure and Witten Laplacian

For a finite set $\Lambda \subset \mathbb{Z}^d$ and a boundary condition η , we define a Hamiltonian by

$$\Phi_{\Lambda,\eta}(x) = \sum_{\substack{i,j\in\Lambda\\i\sim j}} \mathscr{J}(x^i - x^j)^2 + \sum_{i\in Z^d} U(x^i) + 2\sum_{\substack{i\in\Lambda, j\in\Lambda^c\\i\sim j}} \mathscr{J}(x^i - \eta^j)^2$$
(2)

and define a measure on \mathbb{R}^{Λ} by

$$\nu_{\Lambda,\eta} = Z^{-1} e^{-2\Phi_{\Lambda,\eta}(x)} dx_{\Lambda}.$$
(3)

By a standard argument we can define the exterior differentiation d and its adjoint operator d^* with respect to the measure $\nu_{\Lambda,\eta}$.

We call the operator $dd^* + d^*d$ as the Witten Laplacian acting on differential forms.

Vanishing theorem and Hodge-Kodaira decomposition

Suppose that the potential U is decomposed as U = V + W so that V is convex and W is bounded. We denote the supremum and the infimum of W by W_{sup} and W_{inf} , respectively. Then we have

Theorem 1. Suppose $V'' \ge c > 0$, and $2(c+8d \mathscr{J})e^{-2(W_{\sup}-W_{\inf})} > 16d \mathscr{J}$. Then the lowest eigenvalue of $dd^* + d^*d$ on *p*-forms is not less than $\{2(c+8d \mathscr{J})e^{-2(W_{\sup}-W_{\inf})} - 16d \mathscr{J}\}p$. Therefore there is no harmonic *p*-forms $(p \ge 1)$.

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If U has a special form $U(t) = at^4 - bt^2$, we can give different kind of criterion:

Theorem 2. If $\sqrt{3a} - b - 4d \mathscr{J} > 0$, then the lowest eigenvalue of $dd^* + d^*d$ on *p*-forms is not less than $2(\sqrt{3a} - b - 4d \mathscr{J})p$. Therefore there is no harmonic *p*-forms $(p \ge 1)$.

Using these theorems, we can show the Hodge-Kodaira decomposition theorem.

Theorem 3. We have the following : for p = 0,

$$L^{2}(\nu_{\Lambda,\eta}) = \{ \text{ constant functions } \} \oplus \operatorname{Ran}(d^{*}),$$
(4)

and for $p \ge 1$,

$$L^{2}(\nu_{\Lambda,\eta}; \bigwedge^{p}(\mathbb{R}^{\Lambda})^{*}) = \operatorname{Ran}(d) \oplus \operatorname{Ran}(d^{*})$$
(5)

References

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