

Uniqueness of Gibbs measures on $C(\mathbb{R} \rightarrow \mathbb{R})$

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1. h -transformation

We consider the following operator on a domain $D \subseteq \mathbb{R}^d$:

$$L = \frac{1}{2}a_{ij}\partial_i\partial_j + b^i\partial_i + V.$$

We set $L_0 = L - V$. For positive function h , the **h -transform** of L is defined by

$$L^h f = \frac{1}{h}L(hf).$$

More explicitly,

$$L^h f = L_0 + a \frac{\nabla h}{h} \cdot \nabla + \frac{Lh}{h}.$$

If h is a harmonic function, i.e., $Lh = 0$, then L^h has no 0-th order term.

In the sequel, we assume that the semigroup generated by L has the

transition measure: denoting $T_t = e^{tL}$

$$T_t f = \int_D p(t, x, dy) f(y).$$

Then, the transition measure of L^h is given by

$$p^h(t, x, dy) = \frac{1}{h(x)} p(t, x, dy) h(y).$$

2. Invariant function

$p(t, x, dy)$: a transition measure

φ is called a **invariant function** if

$$\varphi(x) = \int_D \varphi(y) p(t, x, dy), \quad \forall t \geq 0.$$

It is easy to see

φ is invariant \Leftrightarrow h -transform by φ is conservative.

principal eigenvalue

For any transition measure $p(t, x, dy)$ associated with L , there exist λ_c so that $L - \lambda$ is subcritical for $\lambda > \lambda_c$ and $L - \lambda$ is supercritical for $\lambda < \lambda_c$. Here

subcritical: Green measure exists

supercritical: no Green measure and no positive harmonic function

λ_c is called a (generalized) **principal eigenvalue**.

We will show that any 1-dimensional (minimal) diffusion process with $\lambda_c = 0$ has an invariant function.

3. One dimensional diffusion processes

$$D = (l_-, l_+).$$

$\{(X_t), P_x\}$: a (minimal) diffusion on D (Dirichlet boundary condition)

$s(x)$: the **scale function**

$dm(x)$: the **speed measure** (standard measure)

ζ : the explosion time

$\frac{d}{dm} \frac{d}{ds}$: the **generator**

Dirichlet form:
$$\mathcal{E}(f, g) = \int_D \frac{df}{ds} \frac{dg}{ds} ds$$

From dm , we define a right continuous non-decreasing function m as

$$m(y) - m(x) = \int_{(x,y]} dm$$

Take any $a \in (l_-, l_+)$ and define

$$S(x) = \int_{(a,x]} \{m(y) - m(a)\} ds(y) = \int_{(a,x]} \{s(x) - s(u)\} dm(u),$$

$$M(x) = \int_{(a,x]} \{s(y) - s(a)\} dm(y) = \int_{(a,x]} \{m(x) - m(u)\} ds(u).$$

- $S(l_+) < \infty \Rightarrow l_+$ is called **exit**.
- $S(l_+) = \infty \Rightarrow l_+$ is called **non-exit**.
- $M(l_+) < \infty \Rightarrow l_+$ is called **entrance**.
- $M(l_+) = \infty \Rightarrow l_+$ is called **non-entrance**.

Feller's criterion:

(X_t) is conservative $\Leftrightarrow S(l_+) = \infty$ and $S(l_-) = \infty$

h -transformation

Let v be a λ -harmonic function, i.e.,

$$\frac{d}{dm} \frac{d}{ds} v = \lambda v.$$

Define $d\hat{m} = v^2 dm$, $d\hat{s} = \frac{ds}{v^2}$. Then

$$(3.1) \quad \frac{1}{v} \left(\frac{d}{dm} \frac{d}{ds} - \lambda \right) (vf) = \frac{d}{d\hat{m}} \frac{d}{d\hat{s}} f.$$

$\frac{d}{d\hat{m}} \frac{d}{d\hat{s}}$ is the h -transform of $\frac{d}{dm} \frac{d}{ds} - \lambda$.

For 1-dimensional diffusions, we have $\lambda_c = \inf \sigma\left(-\frac{d}{dm} \frac{d}{ds}\right)$

Theorem 3.1. Let (X_t) be a diffusion process on D with $\lambda_c = 0$. Then there exist an invariant function.

	left	right	D	eigenvalue	h -transform
case 1	exit ←	exit →	$(0, l)$	$\lambda_0 > 0$	$\varphi_0(x)$
case 2	exit ←	non-exit →/→ ←/← non-entrance	$\begin{cases} (0, \infty) \\ (0, l) \end{cases}$	$\lambda_0 \geq 0$	$s(x) = x$
case 3	exit ←	non-exit →/→ ← entrance	$(0, \infty)$	$\lambda_0 > 0$	$\varphi_0(x)$

4. Gibbs measure on $C(\mathbb{R} \rightarrow \mathbb{R})$

We are given a potential function

- $V : \mathbb{R} \mapsto \mathbb{R} : \text{continuous and non-negative.}$

A Gibbs measure associated with V is formally expressed as

$$\begin{aligned} & \mu(dx) \\ &= Z^{-1} \exp \left\{ -\frac{1}{2} \int_{-\infty}^{\infty} |\dot{x}(t)|^2 dt - \int_{-\infty}^{\infty} V(x(t)) dt \right\} \prod_{t \in \mathbb{R}} dx(t). \end{aligned}$$

Dobrushin-Lanford-Ruelle equation

Precise characterization of Gibbs measure is given as follows. For $I \subseteq \mathbb{R}$, we set $\mathcal{F}_I = \sigma\{x(t); t \in I\}$. Let $P_{s,x}^{t,y}$ be the pinned Brownian motion with $x(s) = x$ and $x(t) = y$. Then a probability measure μ is called a **Gibbs measure** if it satisfies

$$\mu(\cdot | \mathcal{F}_{[s,t]^c})(x(\cdot)) = Z^{-1} \exp\left\{-\int_s^t V(x(u)) du\right\} P_{s,x}^{t,y} \otimes \delta_{x_{[s,t]^c}}.$$

Here Z is a normalizing constant.

We restrict ourselves to the following Gibbs measures:

$$\mathcal{G} = \{\mu \text{ satisfies DLR equation and the family } \{\mu \circ x(t)^{-1}\} \text{ is tight}\}.$$

Schrödinger operator

- $H = \frac{1}{2}\Delta - V$
- $\lambda_0 = \inf \sigma(-H)$: (generalized) principal eigenvalue
- $h = e^{-U}$: (generalized) eigenfunction for λ_0
- If $h \in L^2(\mathbb{R})$, then λ_0 is an eigenvalue.

$$Hh = -\lambda_0 h,$$

$$(H + \lambda_0)^h f = \frac{1}{h}(H + \lambda_0)(hf),$$

$$(H + \lambda_0)^h = \frac{1}{2}\Delta - \nabla U \cdot \nabla.$$

- $((H + \lambda_0)^h, L^2(h^2 dx))$: a self-adjoint operator

- an associated Dirichlet form:

$$\mathcal{E}(f, g) = \frac{1}{2} \int_R (\nabla f, \nabla g) h^2 dx$$

U and V satisfy the following relation:

$$\Delta U - |\nabla U|^2 = 2(\lambda_0 - V).$$

The diffusion operator $\frac{1}{2}\Delta - \nabla U \cdot \nabla$ may generate an explosive diffusion. But we can assume that the associated diffusion is **conservative** by changing an eigenfunction if necessary (Theorem 3.1).

Transition measure associated with $\frac{1}{2}\Delta - V$ is given by

$$p(t, x, y) = E_{0,x}^{t,y}[\exp\{-\int_0^t V(x(s)) ds\}]g(t, x, y)$$

where $E_{0,x}^{t,y}$ stands for the integral with respect to the pinned Wiener measure $E_{0,x}^{t,y}$, and $g(t, x, y)$ is the Gauss kernel

$$g(t, x, y) = \frac{1}{\sqrt{2\pi t}} \exp\left\{-\frac{(y-x)^2}{2}\right\}.$$

Set

$$q(t, x, y) = h(x)^{-1}p(t, x, y)e^{-t\lambda_0}h(y).$$

- $q(t, x, y)dy$: transition measure of the semigroup generated by $(H + \lambda_0)^h$
- h is an invariant function of $H + \lambda_0$

- $h^2 dx$ is an invariant measure

Set

$$\hat{q}(t, x, y) = \frac{q(t, x, y)}{h(y)^2} = \frac{p(t, x, y)e^{-t\lambda_0}}{h(x)h(y)}$$

and

$$\nu(dx) = e^{-2U(x)} dx.$$

$\hat{q}(t, x, y)$ is a density function with respect to ν .

When $\nu(\mathbb{R}) < \infty$, we assume that ν is normalized as $\nu(\mathbb{R}) = 1$.

Then the following is well-known:

- $\nu(\mathbb{R}) = 1 \implies \lim_{t \rightarrow \infty} \sup_{|x|, |y| \leq R} |\hat{q}(t, x, y) - 1| = 0.$
- $\nu(\mathbb{R}) = \infty \implies \lim_{t \rightarrow \infty} \sup_{|x|, |y| \leq R} |\hat{q}(t, x, y)| = 0.$

Theorem 4.1. If λ_0 is an eigenvalue then $\sharp(\mathcal{G}) = 1$, i.e., the **uniqueness** holds.

Sketch of proof

Under the assumption, we have $\nu(\mathbb{R}) = 1$.

- the law of $x(0)$ is ν .

∴ The density function of law of $x(0)$ is

$$\begin{aligned} & \int_{\mathbb{R}^2} \frac{\hat{q}(t, x, z)h(z)^2 \hat{q}(t, z, y)h(y)^2}{\hat{q}(2t, x, y)h(y)^2} \mu^{(x(-t), x(t))}(dx, dy) \\ &= \int_{\mathbb{R}^2} \frac{\hat{q}(t, x, z)h(z)^2 \hat{q}(t, z, y)}{\hat{q}(2t, x, y)} \mu^{(x(-t), x(t))}(dx, dy). \end{aligned}$$

Non-existence

If $V = 0$, then $\sharp(\mathcal{G}) = 0$.

$$\begin{aligned} & \frac{g(t, x, z)g(t, z, y)}{g(2t, x, y)} \mu^{(x(-t), x(t))} (dx, dy) \\ &= \frac{1}{2\pi t} \exp\left\{-\frac{1}{2t}|z - x|^2 - \frac{1}{2t}|y - z|^2\right\} \sqrt{4\pi t} \exp\left\{\frac{1}{4t}|y - x|^2\right\} \\ & \quad \mu^{(x(-t), x(t))} (dx, dy) \\ &= \frac{1}{\sqrt{\pi t}} \exp\left\{-\frac{1}{2t}|z - x|^2 - \frac{1}{2t}|y - z|^2 + \frac{1}{4t}|y - x|^2\right\} \\ & \quad \mu^{(x(-t), x(t))} (dx, dy) \end{aligned}$$